

Spreading in a kinetic reaction-transport equation for population dynamics

CCSAMM

Nils Caillerie

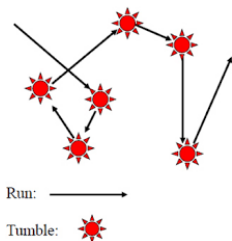
Georgetown University

September 20th 2017

collaboration with E. Bouin (U. Paris-Dauphine)
work performed during my PhD at Université Lyon 1

Spreading in a kinetic reaction-transport equation

Applications: *Escherichia coli*¹, *Rhinella marina*² (cane toad)



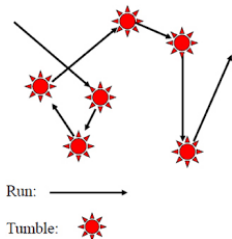
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Run



- Exponential time with mean 1
- admissible velocity set: $V \subset \mathbb{R}^d$ (compact)

Tumble

- velocity redistribution: $M \in L^1(V)$ such that $\int_V vM(v)dv = 0$

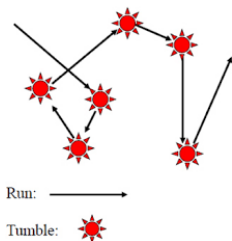
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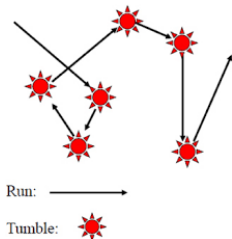
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- velocity redistribution:
 $M \in L^1(V)$ such that
 $\int_V vM(v)dv = 0$

+ **Reproduction** at rate $r > 0$ and intra-specific competition (KPP type)
+ homogeneous environment (No chemotactic effect !!!)

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Spreading in a kinetic reaction-transport equation

Chapman-Kolmogorov equation:

$$\partial_t f + v \cdot \nabla_x f = M(v)\rho - f + r\rho(M(v) - f), \quad (t, x, v) \in \mathbb{R}_+ \times \mathbb{R}^d \times V. \quad (1)$$

Density: $\rho(t, x) := \int_V f(t, x, v) dv$

We will assume $V = B(0, v_{\max})$, $v_{\max} < +\infty$.

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How fast does the population spread in its environment?

Spreading in a kinetic reaction-transport equation

First approach: Diffusion-approximation $(t, x, v) \rightarrow (\frac{t}{\varepsilon^2}, \frac{x}{\varepsilon}, v)$, $r \rightarrow \varepsilon^2 r$

$$\varepsilon^2 \partial_t f^\varepsilon + \varepsilon v \cdot \nabla_x f^\varepsilon = M \rho^\varepsilon - f^\varepsilon + \varepsilon^2 r \rho^\varepsilon (M - f^\varepsilon) \quad (2)$$

Assume $\int_V v M(v) = 0$ and $\int_V |v|^2 M(v) = \theta < +\infty$, then³

$$\lim_{\varepsilon \rightarrow 0} f^\varepsilon(t, x, v) = \rho(t, x) M(v), \quad (3)$$

where $\partial_t \rho - \theta \partial_{xx}^2 \rho = r \rho (1 - \rho)$ (Fisher-KPP).

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From Kolmogorov-Petrovskii-Piskunov ('37):

- travelling waves for all $c \geq 2\sqrt{r\theta}$
- compactly supported initial data: spreading at speed $2\sqrt{r\theta}$

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Spreading in a kinetic reaction-transport equation

This approach may underestimate the speed of propagation, due to:

- chemotactic effect (e.g. *Escherichia coli*⁴)
- strongly biased random walks (e.g. *Rhinella marina*⁵)

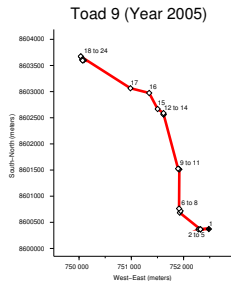
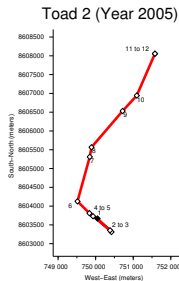
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data: courtesy of G. Brown, B. Phillips, R. Shine

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Spreading in a kinetic reaction-transport equation

Other approaches:

- 1 Travelling wave solutions
- 2 Hyperbolic limits

Travelling wave solutions

Looking for solution f of (1) of the form $f(t, x, v) = h(x \cdot e - ct, v)$:

$$(v \cdot e - c)\partial_1 h = M\rho_h - h + r\rho_h(M - h),$$

with $h(z, v) \rightarrow M(v)$ as $z \rightarrow -\infty$ and $h(z, v) \rightarrow 0$ as $z \rightarrow +\infty$

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(Note: supersolution)

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(Note: supersolution)

2. Exponential decay at the edge of the front: $\tilde{h}(z, v) \simeq e^{-\lambda z} Q(v)$.

Formally, we get the spectral problem: Find $(c\lambda, Q_\lambda)$ such that

$$\lambda c Q_\lambda = (\lambda v \cdot e - 1) Q_\lambda + (r + 1) M \int_V Q_\lambda(v) dv.$$

Travelling wave solutions

Solving the spectral problem gives the dispersion relation

$$(r + 1) \int_V \frac{M(v)}{1 + \lambda(c - v \cdot e)} dv = 1. \quad (4)$$

When, $d = 1$ and $M > 0$,

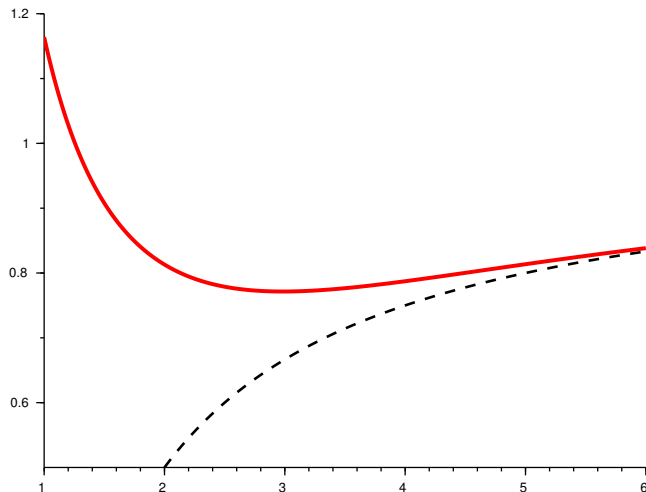
$$\lim_{c \rightarrow v_{\max} - \frac{1}{\lambda}} (r + 1) \int_V \frac{M(v)}{1 + \lambda(c - v \cdot e)} dv = +\infty.$$

Theorem (Bouin-Calvez-Nadin 2015)

Suppose $d = 1$ and $\inf_{v \in V} M(v) > 0$. Then, for $\lambda > 0$, there exists a unique $c(\lambda) > 0$ such that (4) holds. Moreover, there exist travelling wave solutions of (1) for all $c \geq c^* := \min_{\lambda > 0} c(\lambda)$.

Travelling wave solutions

$\lambda \mapsto c(\lambda)$ for $V = [-1, 1]$ and $M \equiv \frac{1}{2}$



Travelling wave solutions

Construction of travelling wave solutions for $c \geq c^*$:

- There exists a solution Q_λ to the spectral problem
- sub- and super-solution using Q_λ
- comparison principle

Non-existence of travelling wave solutions for $c < c^*$:

- relies on the fact that $c'(\lambda) = 0$

Travelling wave solutions

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Then, solution of the spectral problem is given $(v_{\max} - 1/\lambda, \mu)$, where

$$\mu = \left(1 - \frac{(r + 1)}{\lambda} \int_V \frac{M(v)}{v_{\max} - v' \cdot e} dv' \right) \delta_{v_{\max} e} + \frac{dv}{\lambda(v_{\max} - v \cdot e)}$$

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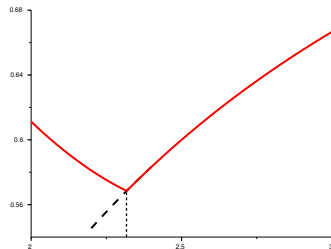
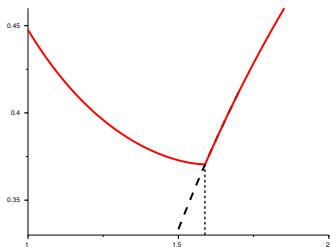
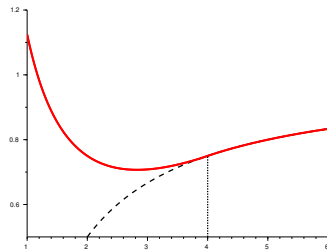
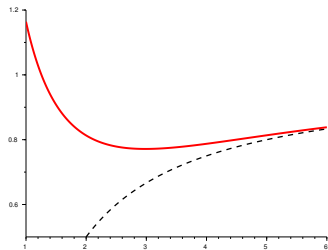
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Occurs in the most simple cases: for $d = 2$, $V = D(0, 1)$ and $M \equiv \frac{1}{\pi}$,

$$\frac{r+1}{\lambda} \int_V \frac{M(v)}{v_{\max} - v \cdot e} dv = \frac{r+1}{2\lambda}.$$

Travelling wave solutions



Travelling wave solutions

Result for general d and possibly vanishing M :

Theorem (Bouin, NC, submitted (2017))

Under previous assumptions, there exist travelling wave solutions for all $c \geq c^ = \min_{\lambda > 0} c(\lambda)$.*

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Existence for $c \geq c^*$: exactly as in Bouin Calvez & Nadin's paper (sub- and super-solution)

Non-existence when $c < c^*$: requires more work (Hyperbolic limits)

Hyperbolic limits

Study (1) in the hyperbolic scale: $(t, x, v) \rightarrow (\frac{t}{\varepsilon}, \frac{x}{\varepsilon}, v)$.

$$\partial_t f^\varepsilon + v \cdot \nabla_x f^\varepsilon = \frac{1}{\varepsilon}(M\rho^\varepsilon - f^\varepsilon) + \frac{r}{\varepsilon}\rho^\varepsilon(M - f^\varepsilon).$$

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WKB Ansatz: $\varphi^\varepsilon := -\varepsilon \ln(\frac{f^\varepsilon}{M})$, equivalently $f^\varepsilon = M e^{-\frac{\varphi^\varepsilon}{\varepsilon}}$.

$$\partial_t \varphi^\varepsilon + v \cdot \nabla_x \varphi^\varepsilon + r = (1 + r) \int_V M(v') (1 - e^{\frac{\varphi^\varepsilon - \varphi'^\varepsilon}{\varepsilon}}) dv' + r\rho^\varepsilon$$

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Suppose $\varphi^\varepsilon \rightarrow \varphi^0$, then (formally)

$$f^\varepsilon \rightarrow 0 \quad \text{on } \{\varphi^0 > 0\},$$

$$f^\varepsilon \rightarrow M \quad \text{on } \{\varphi^0 = 0\}.$$

Hyperbolic limits

What is φ^0 ? Thanks to Lipschitz uniform bounds⁶ on φ^ε , φ^0 should be independent of ν .

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Let us assume that $\varphi^\varepsilon = \varphi^0 - \varepsilon \ln(Q) + \mathcal{O}(\varepsilon^2)$. Then, formally,

$$\partial_t \varphi^0 + v \cdot \nabla_x \varphi^0 + r = (1+r) \int_V M(v') \left(1 - \frac{Q(v')}{Q(v)}\right) dv' + r e^{-\frac{\varphi^0}{\varepsilon}} \int_V Q(v') dv'$$

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Let $p := \nabla_x \varphi^0$ and $H = r - \partial_t \varphi^0$. On $\{\varphi^0 > 0\}$,

$$Q(v) = \frac{1+r}{1+H-v \cdot p} \int_V M(v') Q(v') dv',$$

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Hence,

$$(1+r) \int_{\mathcal{V}} \frac{M(v)}{1 + H(p) - v \cdot p} dv = 1. \quad (5)$$

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Remark: $c(\lambda) = \frac{H(\lambda e)}{\lambda}$

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Hence,

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Remark: $c(\lambda) = \frac{H(\lambda e)}{\lambda}$

Theorem (Bouin 2015)

For $d = 1$, $M > 0$ and for $\varphi(0, x, v) = \varphi_0(x)$, the sequence $(\varphi^\varepsilon)_\varepsilon$ converges uniformly locally to φ^0 which is the viscosity solution of

$$\begin{cases} \min(\partial_t \varphi^0 + H(\nabla_x \varphi^0) + r, \varphi^0) = 0, \\ \varphi^0(0, \cdot) = \varphi_0, \end{cases}$$

where H is implicitly defined by (5).

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where H is implicitly defined by (6) when such H exists, and $H(p) = v_{\max} |p| - 1$ otherwise.

φ^0 is a viscosity super-solution : let $\psi \in C^1(\mathbb{R}_+ \times \mathbb{R}^d)$ such that $\varphi^0 - \psi$ has a global strict minimum at $(t^0, x^0) \in \mathbb{R}_+^* \times \mathbb{R}^d$.

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$\varphi^\varepsilon - \psi^\varepsilon$ has a minimum at $(t^\varepsilon, x^\varepsilon, v^\varepsilon)$.

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At $(t^\varepsilon, x^\varepsilon, v^\varepsilon)$,

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$$\begin{aligned} \frac{\partial \psi^\varepsilon}{\partial t} + v^\varepsilon \cdot \nabla_x \psi^\varepsilon + r &= (1+r) \left(1 - \int_V M' e^{\frac{\varphi^\varepsilon - \varphi'^\varepsilon}{\varepsilon}} dv' \right) + r \int_V e^{-\frac{\varphi'^\varepsilon}{\varepsilon}} dv' \\ &\geq (1+r) \left(1 - \int_V M' e^{\frac{\psi^\varepsilon - \psi'^\varepsilon}{\varepsilon}} dv' \right) + r \int_V e^{-\frac{\varphi'^\varepsilon}{\varepsilon}} dv' \\ &= (1+r) \left(1 - \frac{1}{Q} \int_V M' Q' dv' \right) + r \int_V e^{-\frac{\varphi'^\varepsilon}{\varepsilon}} dv' \\ &= (1+r) \left(1 - \frac{1}{(r+1)Q(v^\varepsilon)} \right) + r \int_V e^{-\frac{\varphi'^\varepsilon}{\varepsilon}} dv' \\ &= -H(p^0) + v^\varepsilon \cdot p^0 + r \int_V e^{-\frac{\varphi'^\varepsilon}{\varepsilon}} dv' \end{aligned}$$

1st case: Then, $(1+r) \int_V M(v)Q(v)dv = 1$ and $Q \in L^\infty(V)$

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Take the limit $\varepsilon \rightarrow 0$:

$$\frac{\partial \psi}{\partial t}(t^0, x^0) + v^* \cdot \nabla_x \psi(t^0, x^0) + r \geq -H(\nabla_x \psi(t^0, x^0)) + v^* \cdot \nabla_x \psi(t^0, x^0)$$

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Q_K is bounded: same procedure, then take $K \rightarrow +\infty$

Hyperbolic limits

Back to spreading issues:

Let f be a travelling wave solution: $f(t, x, v) = h(x \cdot e^{-ct}, v)$ and $f^\varepsilon = h\left(\frac{x \cdot e^{-ct}}{\varepsilon}, v\right)$.

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- 2 φ^0 solves
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- 3 Hopf-Lax formula:
$$\varphi^0(t, x) = \max\left(\min_{y \in \mathbb{R}^d} \left\{ tL\left(\frac{x-y}{t}\right) + \varphi^0(0, y) \right\}, 0\right),$$
 where L is the Legendre transform of H : $L(p) = \sup_{q \in \mathbb{R}^d} \{p \cdot q - H(q)\}$

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- 4 $\varphi^0 = 0 \implies L\left(\frac{x}{t}\right) \leq 0$ and $x \cdot e \leq \frac{H(\lambda e)}{\lambda} t = c(\lambda)t$ for $\lambda > 0$.
- 5 Since $h(z, v) \xrightarrow{z \rightarrow -\infty} M(v)$ and $h(z, v) \xrightarrow{z \rightarrow +\infty} 0$, we have $c \geq c^*$.

- $r = 0$: Bouin-Calvez 2012 and NC 2017
- Unbounded velocity set (superlinear spreading): Bouin-Calvez-Grenier-Nadin 2016 (submitted)
- $r = 0$ and force terme: NC (work in progress)
- More general reaction terms (In 1D !!!): Bouin 2016.
- genetic trait structured population: Bouin-Mirrahimi (2015)

Thank you for your attention !

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